

Energy interactions in homogeneously sheared magnetohydrodynamic flows

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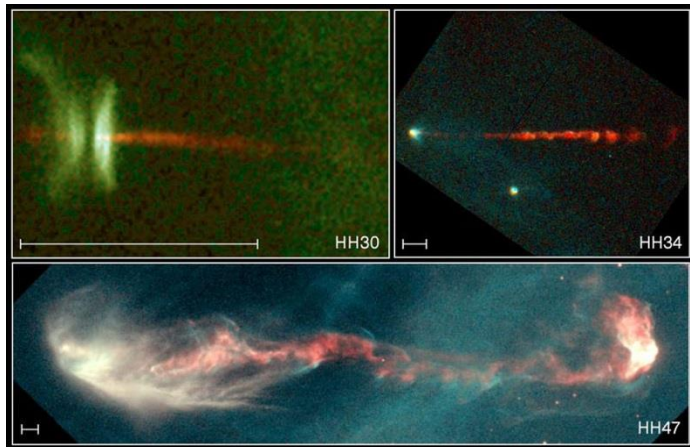
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Outline

- ❖ Motivation
- ❖ Objectives
- ❖ Approach
- ❖ Simulation setup
- ❖ Results
- ❖ Conclusions

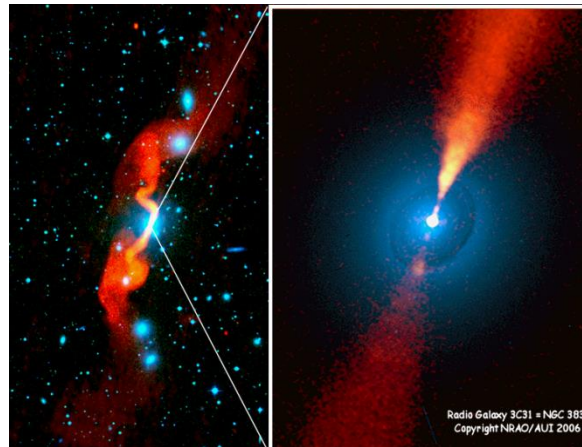
Motivation

- ❖ Plasma is ubiquitous – constitutes $\sim 99\%$ of baryonic matter
 - Widely observed in nature and engineering
- ❖ Magnetic field is generally known to stabilize instabilities
 - Kelvin-Helmholtz or Richtmeyer-Meskov instability



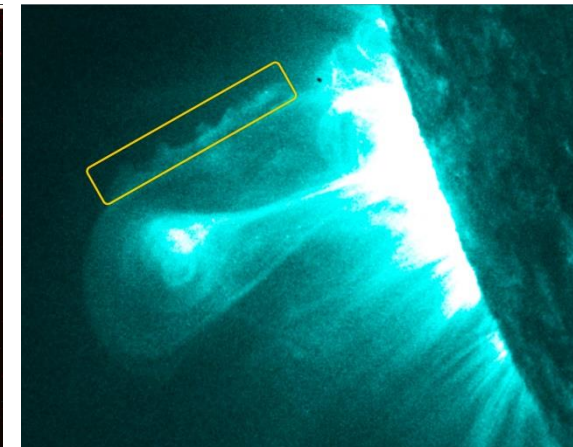
Jets from young stars

Photo credit: C. Burrows (ST Scl), J. Hester (ASU), J. Morse (ST Scl), NASA



Active galactic nuclei of radio galaxy 3C31

Photo credit: NRAO/AUI 2006



Kelvin-Helmholtz instability in a coronal mass ejection

Photo credit: NASA AIA/SDO

Objectives

- ❖ Investigate the evolution of velocity and magnetic fields as a function of
 - Magnetic field strength (B_0)
 - Perturbation wavevector orientation (β)
- ❖ Investigate energy exchange between perturbation velocity and magnetic fields

Theory

- ❖ Considering perturbations in shearnormal plane:

$$\kappa_1(t) = \kappa_0 \cos \beta, \kappa_2(t) = -\kappa_0 \cos \beta St, \\ \kappa_3(t) = \kappa_0 \sin \beta$$

κ_0 is the initial wavevector magnitude

- ❖ Dimensionless quantities using reference distance, time, velocity:

$$\kappa_{ref} = \kappa_0; t_{ref} = \frac{1}{S}; u_{ref} = V_A = \frac{B_0}{\sqrt{\rho \mu_0}}$$

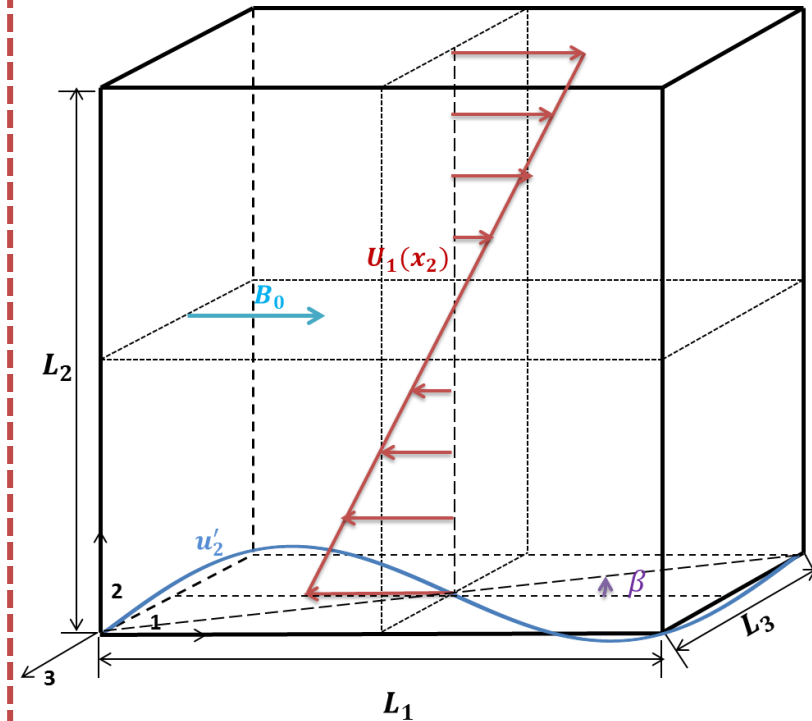
- ❖ Dimensionless governing equations in spectral space:

$$\frac{d\hat{u}_i^*}{d\tau} = \hat{u}_2^* \left(-\delta_{i1} + \frac{2\cos\beta\kappa_i^*}{1 + \cos^2\beta\tau^2} \right) + iR_A^* \hat{B}_i^*$$

$$\frac{d\hat{B}_i^*}{d\tau} = \hat{B}_2^* \delta_{i1} + iR_A^* \hat{u}_i^*$$

$$R_A^* = \frac{\kappa_0 V_A \cos \beta}{S} = R_A \cos \beta$$

Where, $i = 1 - 3$



Problem Setup

❖ The evolution of velocity and magnetic fields decided by R_A^*

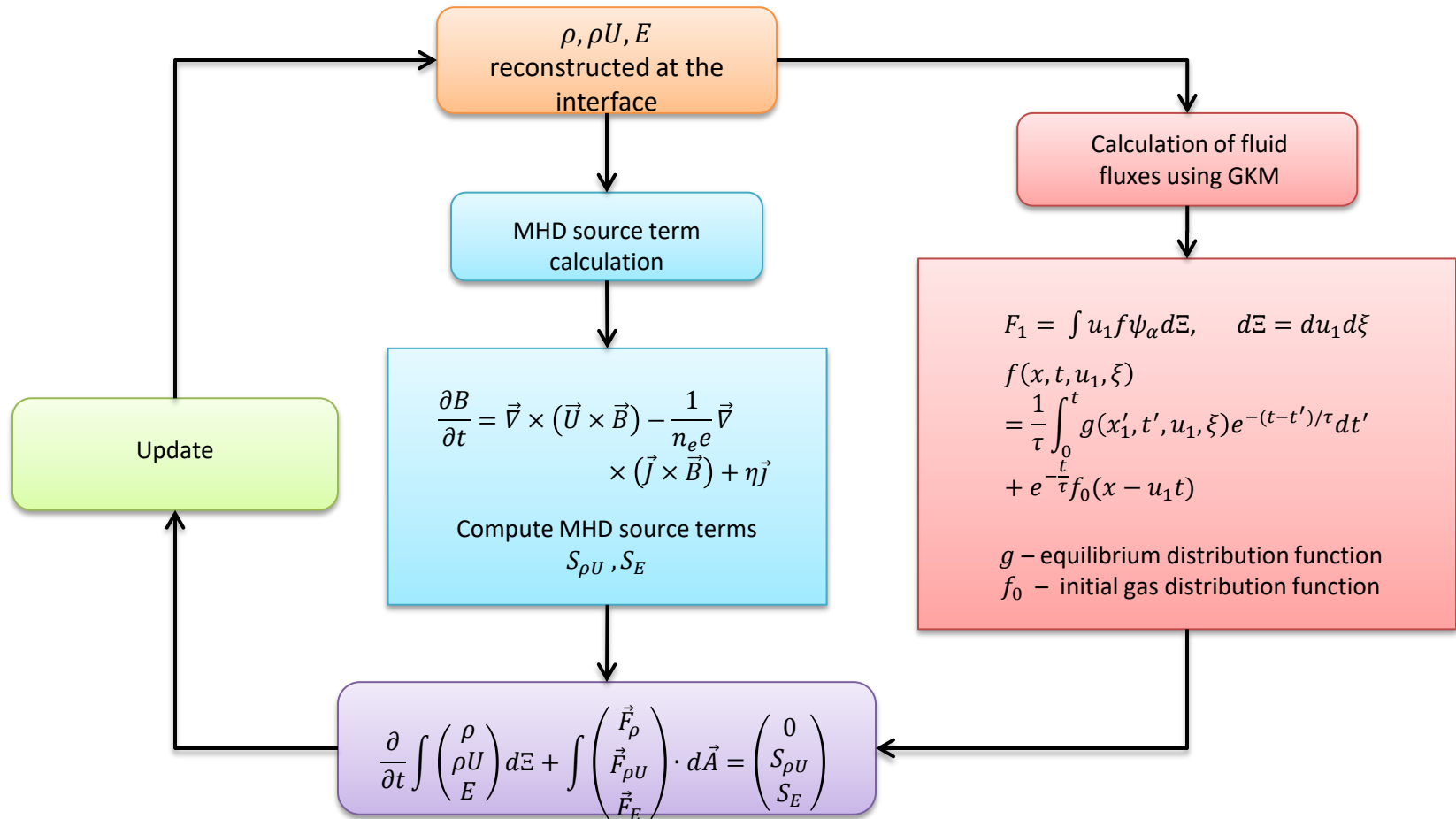
- R_A^* : Ratio of shear ($\tau_s = \frac{1}{s}$) and magnetic time scales ($\tau_B = \frac{1}{\kappa_0 V_A \cos \beta}$)
- Magnetic frequency: $\frac{1}{\tau_B} \propto \cos \beta$

❖ For a given V_A, S, κ_0 :

- Streamwise perturbations ($\beta = 0^\circ$): R_A^* is maximum
 - Highest harmonic exchange between \hat{u}_i^* and $\hat{B}_i^* \Rightarrow$ equi-partition between kinetic and magnetic energies can be observed
- Spanwise perturbations ($\beta = 90^\circ$): $R_A^* = 0$
 - No exchange between \hat{u}_i^* and $\hat{B}_i^* \Rightarrow$ no evolution of \hat{B}_i^*
 - \hat{u}_i^* equations are pressure-released
 - Kinetic energy grows quadratically unaffected by magnetic field strength
- Intermediate orientations ($\beta = 30^\circ, 60^\circ$):
 - Mixed behavior depending on the orientation

Numerical Scheme

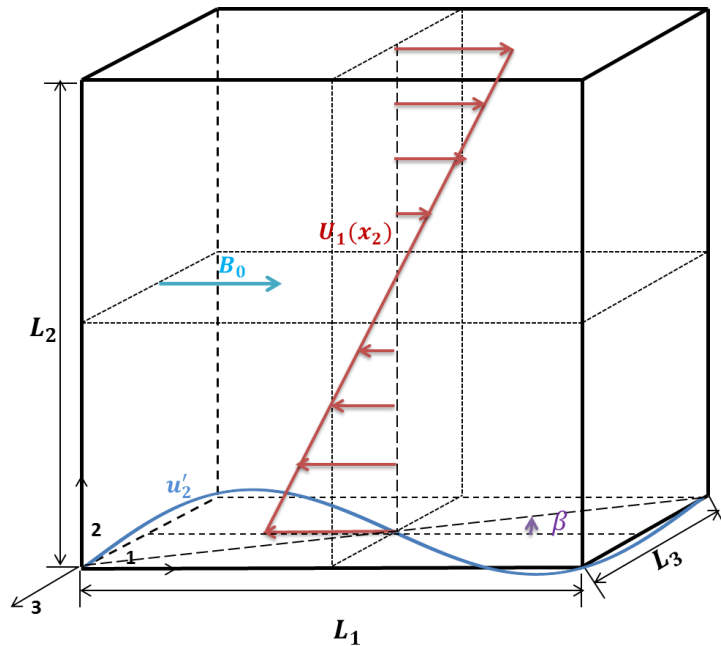
- Magneto-Gas Kinetic Method^{1,2} (MGKM) solves fluid equations with the simplified Boltzmann equation and the magnetic field equations *separately*.



¹Xu, K., Journal of Computational Physics 171, 289-335 (2001)

²Araya, D.B. et al. , ASME, Vol. 137, 081302-(1-11), Aug., 2015

Simulation setup

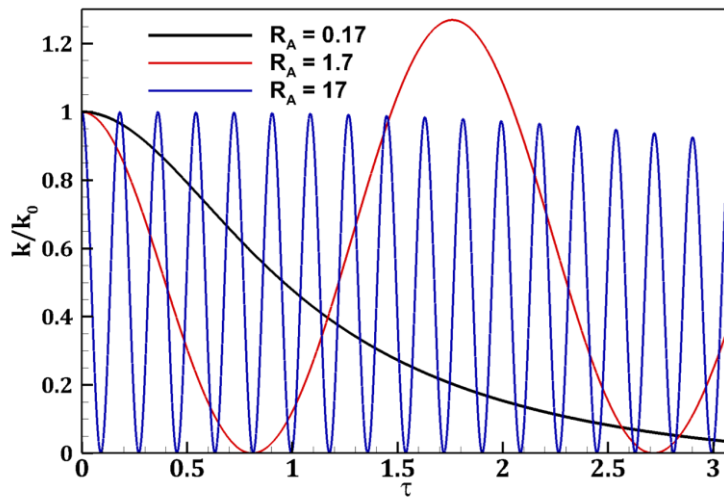


- ❖ MHD homogeneous shear simulation conditions
 - Temperature = 300 K, Density = 1 kg/m³
 - Reynolds number = 770
 - Magnetic Reynolds Number = 195
 - Gradient Mach number, $M_g = 0.03$
- ❖ Boundary conditions:
 - 1 and 3 planes: periodic boundaries
 - 2 plane: shear periodic boundary

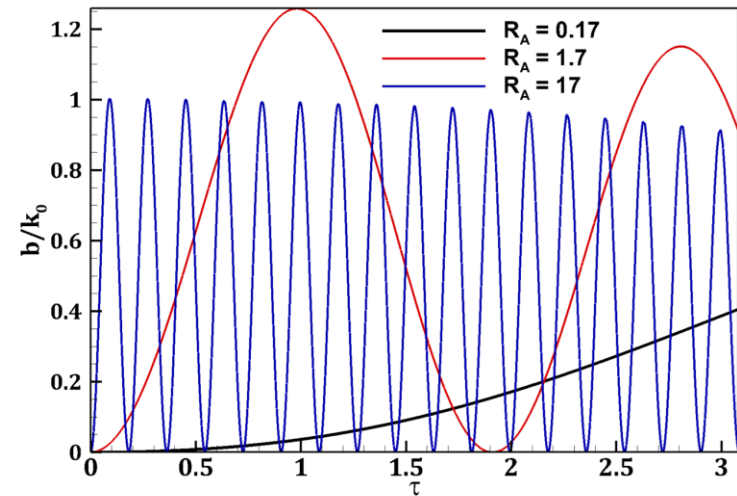
| Magnetic field strength(B_0) | $V_A = \frac{B_0}{\sqrt{\mu_0 \rho}}$ | Perturbation orientation (β) | $R_A = \frac{V_A \kappa_0}{S}$ | $R_A^*(\beta = 0^\circ, 30^\circ, 60^\circ, 90^\circ)$ |
|----------------------------------|---------------------------------------|--------------------------------------|--------------------------------|--|
| 0.0003 T | 0.26m/s | 0°, 30°, 60°, 90° | 0.17 | 0.17, 0.14, 0.085, 0 |
| 0.003 T | 2.6m/s | 0°, 30°, 60°, 90° | 1.7 | 1.7, 1.4, 0.85, 0 |
| 0.03 T | 26m/s | 0°, 30°, 60°, 90° | 17 | 17, 14, 8.5, 0 |

Simulation parameters

Effect of magnetic field strength at $\beta = 0^\circ$



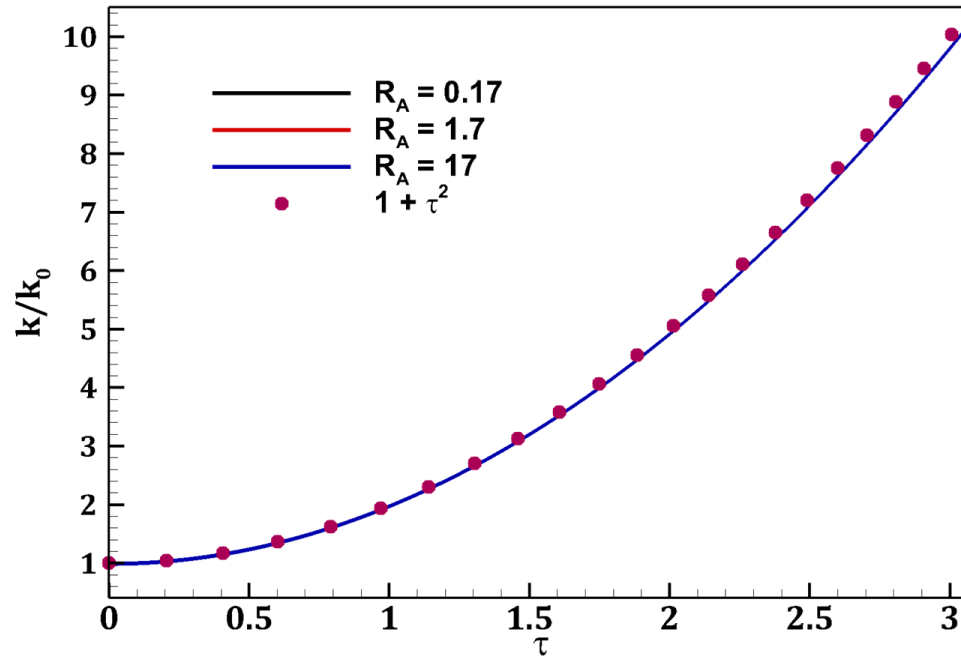
Evolution of perturbation kinetic energy (k)
with respect to initial kinetic energy (k_0)



Evolution of perturbation magnetic energy (b)
with respect to initial kinetic energy (k_0)

- ❖ As R_A increases, wave-like behavior increases
- ❖ $R_A = 0.17$:
 - Kinetic energy monotonically decays to zero \Rightarrow Transfer to mean kinetic energy via the action of pressure
 - Magnetic energy grows monotonically
- ❖ $R_A = 1.7$:
 - Solution exhibits wave-like behavior and overshoots the initial kinetic energy \Rightarrow Conversion to perturbation kinetic energy from mean energy
- ❖ $R_A = 17$, solution oscillates and decays gradually

Effect of magnetic field strength at $\beta = 90^\circ$

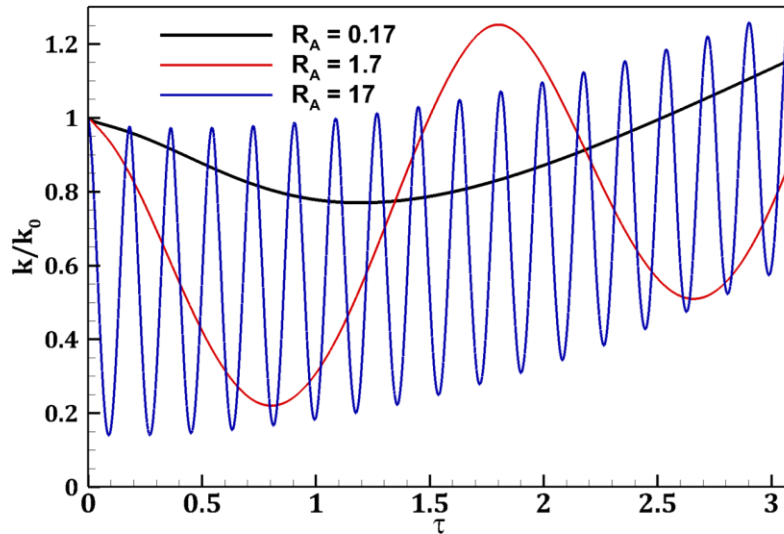


Evolution of perturbation kinetic energy (k) with respect to initial kinetic energy (k_0)

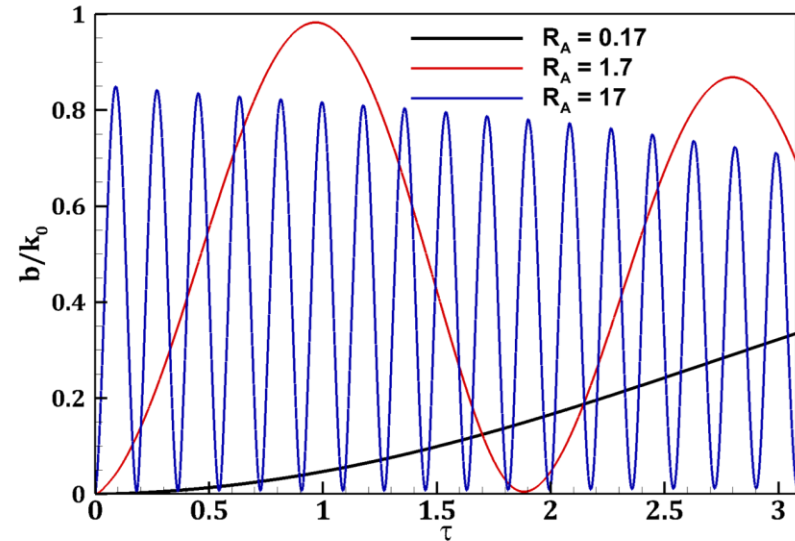
- ❖ Kinetic energy grows quadratically at all values of R_A
- ❖ The evolution matches the pressure-released Burgers evolution given by Simone et al.

$$\frac{k}{k_0} = 1 + \tau^2$$

Effect of magnetic field strength at $\beta = 30^\circ$



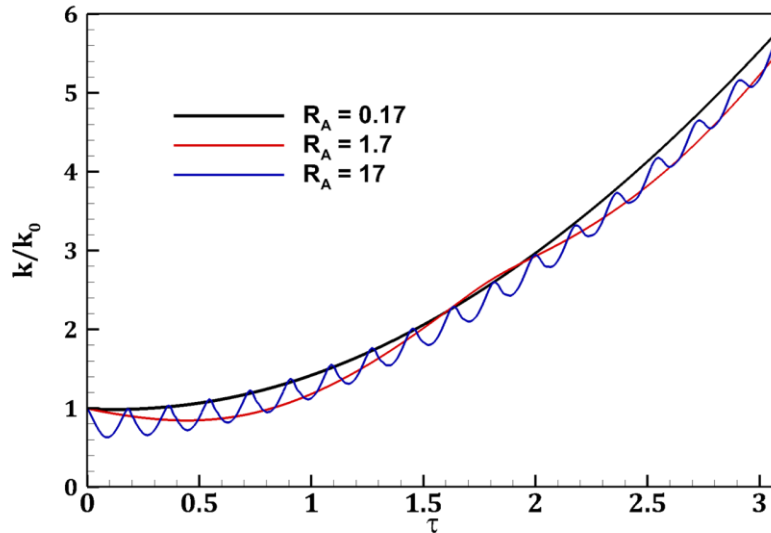
Evolution of perturbation kinetic energy (k)
with respect to initial kinetic energy (k_0)



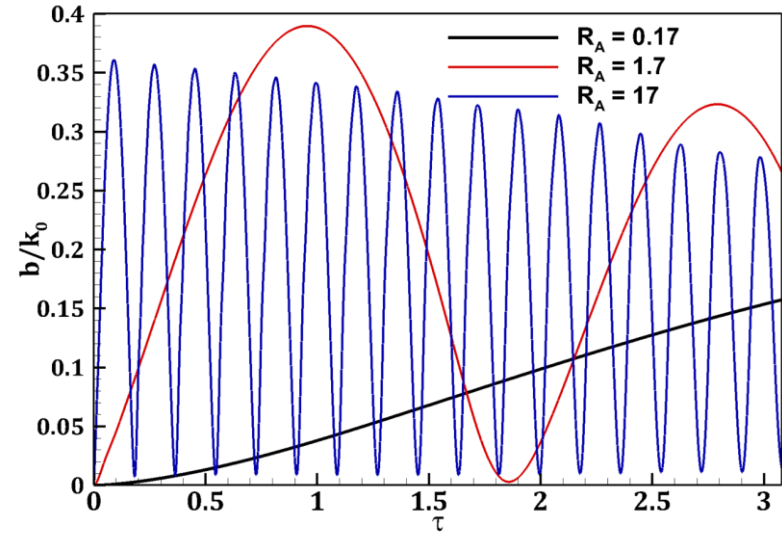
Evolution of perturbation magnetic energy (b)
with respect to initial kinetic energy (k_0)

- ❖ At all R_A , kinetic energy decreases initially, but increases at later times \Rightarrow spanwise behavior dominates at later times
- ❖ Magnetic energy evolution similar to $\beta = 30^\circ$, except at lower magnitudes

Effect of magnetic field strength at $\beta = 60^\circ$



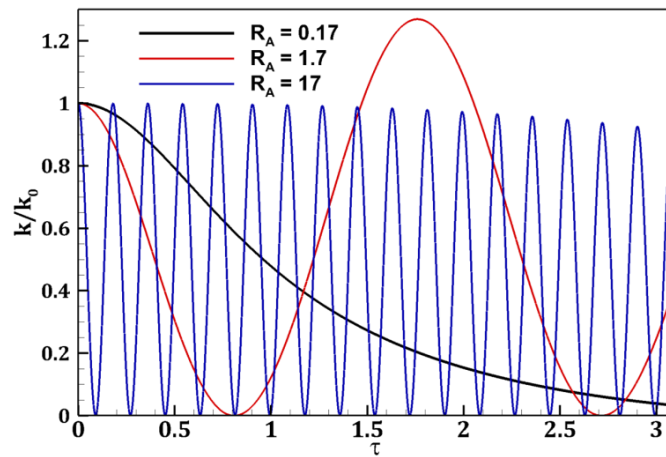
Evolution of perturbation kinetic energy (k)
with respect to initial kinetic energy (k_0)



Evolution of perturbation magnetic energy (b)
with respect to initial kinetic energy (k_0)

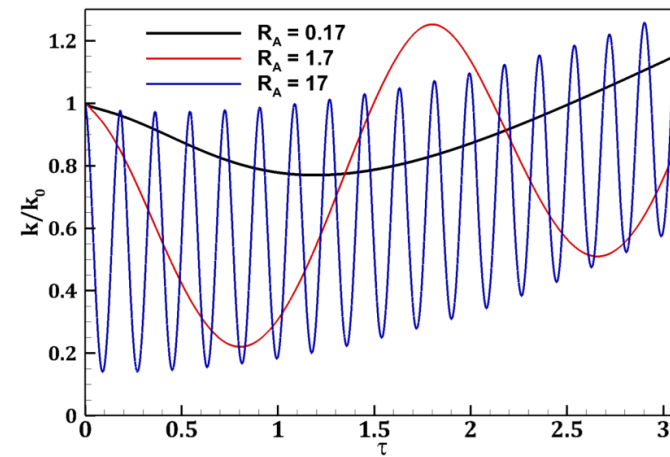
- ❖ At all R_A , kinetic energy exhibits similar growth; oscillations about $R_A = 0.17$ evolution decrease in magnitude
- ❖ Magnetic energy plots similar to $\beta = 0^\circ, 30^\circ$; except at lower magnitude
- ❖ Predominantly spanwise behavior

The effect of orientation (β)

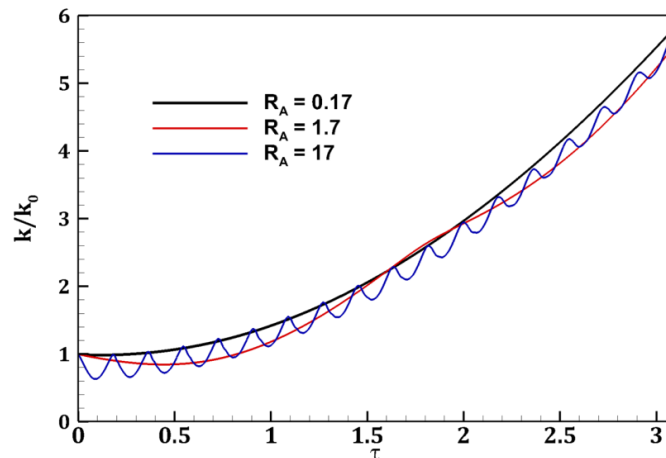


$\beta = 0^\circ$

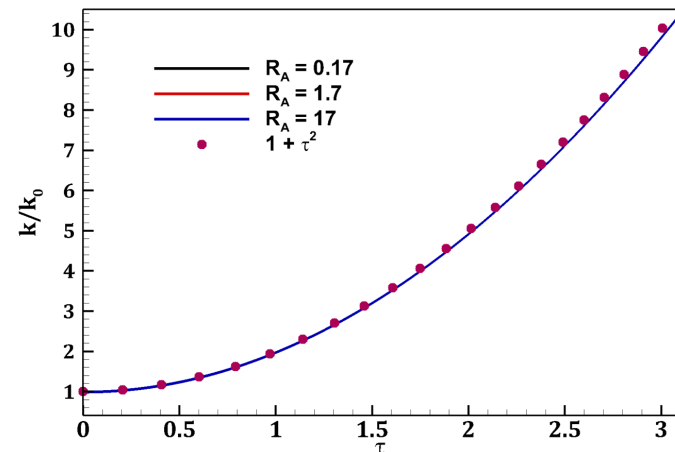
As β increases, the effect of magnetic field decreases



$\beta = 30^\circ$



$\beta = 60^\circ$



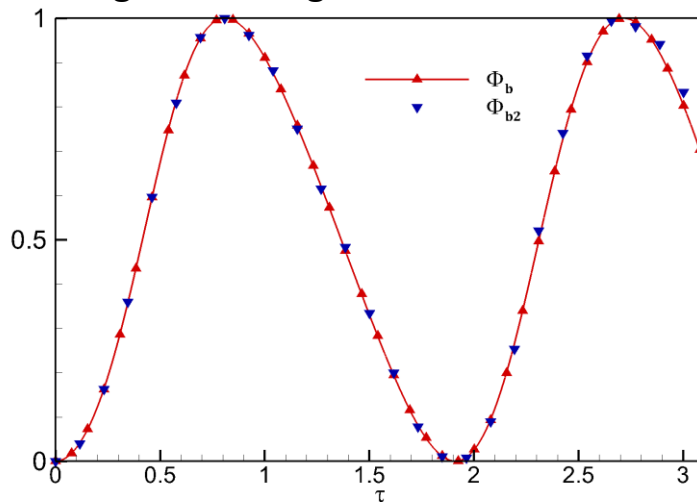
$\beta = 90^\circ$

Harmonic energy exchanges

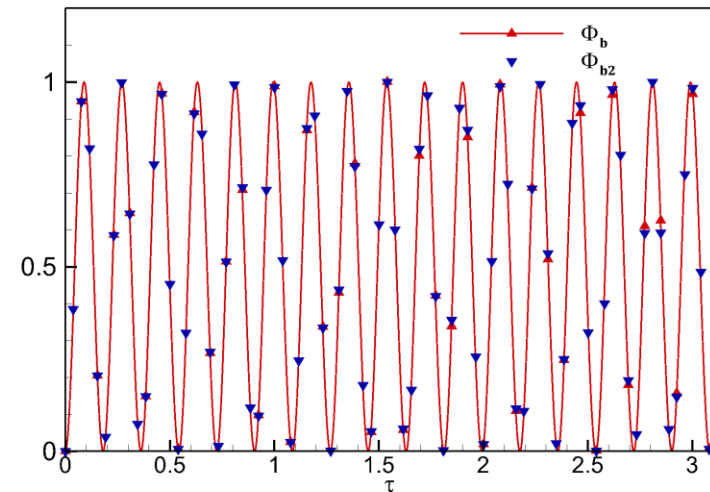
- ❖ Kinetic and magnetic energies exhibit similar magnitudes at $\beta = 0^\circ$
- ❖ Equi-partition potentials defined as:

$$\Phi_b = \frac{b}{k + b}$$
$$\Phi_{b2} = \frac{b_2}{k_2 + b_2}$$

Where, k, k_2 are total, 2-component kinetic energies, and b, b_2 are total, 2-component magnetic energies



$$R_A = 1.7$$



$$R_A = 17$$

- ❖ Oscillates about 0.5 with frequency of π/R_A
 - Equi-partition of kinetic and magnetic energies and their 2-component

- ❖ $R_A^* = \frac{V_A \kappa_0 \cos \beta}{S}$ characterizes the effect of magnetic field strength and wavevector orientation
- ❖ As magnetic field strength increases, R_A^* increases :
 - Wave-like behavior results
- ❖ As β increases from 0° to 90° , R_A^* decreases:
 - $\beta = 0^\circ$: highest harmonic exchange, equi-partition at $R_A = 1.7, 17$
 - $\beta = 90^\circ$: no exchange, pressure-released behavior at all magnetic field strengths
 - As β varies from 0° to 90° , the effect of magnetic field decreases

Thank you!
Questions?